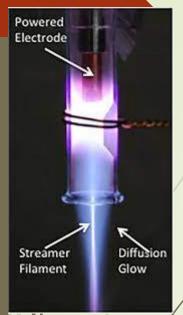
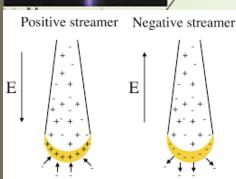
Simulations of positive and negative streamers in the AMReX environment

Ilija Simonović, Danko Bošnjaković and Saša Dujko







What are streamers?

Thin channels of weakly-ionized nonstationary plasma produced by an ionization front that moves through non-ionized matter

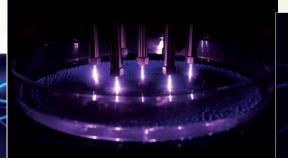
Streamers in nature: lighting and sprite discharges in the upper planetary atmospheres

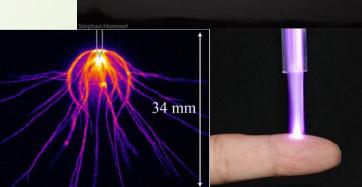
Applications of streamers:

o ignition of high-intensity discharge lamps

treatment of polluted gases and water

o plasma medicine





Streamer modelling

- Motivation:
- better understanding of streamer physics
- o optimization of streamer applications
- Types of streamer models
- o Particle models
- o Fluid models
- Hybrid models

Fluid models of streamers

- Number densities of charged particles are represented by continuous functions
- Time evolution of these number densities are represented by the fluid equations
- Fluid models are generally more computationally efficient than the particle models and are much simpler than the hybrid models

First-order fluid model

Advection diffusion reaction equation for the time evolution of the number density of electrons:

$$\frac{\partial n_e}{\partial t} + \nabla (n_e \mathbf{W} - \mathbf{D} \nabla n_e) = n_e (\alpha - \eta) |\mathbf{W}| + S_{ph}$$

Reaction equations for the time evolution of the number densities of ions:

$$\frac{\partial n_p}{\partial t} = n_e \alpha |\mathbf{W}| + S_{ph} \qquad \frac{\partial n_n}{\partial t} = n_e \eta |\mathbf{W}|$$

Local field approximation

Total electric field:

$$\mathbf{E} = \mathbf{E}_{applied} - \nabla \Phi_{space_charge}$$

$$\Delta \Phi_{space_charge} = -q_e \frac{n_p - n_e - n_n}{\varepsilon_0}$$

Photoionization

- In air photons emitted from excited nitrogen molecules can ionize oxygen molecules
- Zheleznyak model:

$$S_{ph}(\mathbf{r}) = \int d^3r' \frac{I(\mathbf{r}')f(|\mathbf{r}-\mathbf{r}'|)}{4\pi|\mathbf{r}-\mathbf{r}'|^2}$$

■ The photon source term

$$I(\mathbf{r}) = \frac{p_q}{p + p_q} \xi^B \frac{\nu_u}{\nu_i} S_i(\mathbf{r})$$

Ionization source term:

$$S_i(\mathbf{r}) = \alpha |\mathbf{W}|$$

■ The absorption function:

$$\int_0^\infty f(r)dr = 1$$

Photoionization

The absorption function can be represented as:

$$f(|\mathbf{r} - \mathbf{r}'|) = p_{O_2}^2 |\mathbf{r} - \mathbf{r}'| \sum_{j=1}^N A_j^B e^{-\lambda_j^B p_{O_2} |\mathbf{r} - \mathbf{r}'|}$$

Leading to the decomposition of the source term:

$$S_{ph} = \sum_{j=1}^{N} S_{ph,j}$$

$$S_{ph}(\mathbf{r}) = \frac{p_q}{p + p_q} \xi^B \frac{\nu_u}{\nu_i} \sum_{j=1}^{N} p_{O_2}^2 A_j^B \int d^3r' \frac{S_i(\mathbf{r}') e^{-\lambda_j^B p_{O_2}^2 |\mathbf{r} - \mathbf{r}'|}}{4\pi |\mathbf{r} - \mathbf{r}'|}$$

■ The components of the source term can be determined by solving a set of Helmholtz equations:

$$\left(\Delta - (p_{O_2}\lambda_j^B)^2\right)S_{ph,j} = -\left(A_j^B p_{O_2}^2 \frac{p_q}{p + p_q} \xi^B \frac{\nu_u}{\nu_i}\right)S_i$$

Parametrizations: Luque, Bourdon 2 term, Bourdon 3 term

Spatial discretization: Finite volume method

- Scalar variables are defined at cell centers
- Vector variables are defined at cell faces
- Electric field components at the cell centers are interpolated from the cell faces
- Interpolation of the number density of electrons from the cell centers to the cell faces, to calculate electron fluxes, is not trivial!
- First order upwind scheme creates too much numeric diffusion, while linear interpolation (central differencing) creates strong numeric oscillations.
- For this reason, flux limiting schemes are used.

Koren flux limiter

The flux of electrons is calculated as:

$$f_{i+\frac{1}{2}} = \begin{cases} W_{i+\frac{1}{2}} \left(n_i + \frac{1}{2} \phi \left(r_{i+\frac{1}{2}}^+ \right) (n_i - n_{i-1}) \right), & \text{for } W_{i+\frac{1}{2}} > 0 \\ W_{i+\frac{1}{2}} \left(n_{i+1} + \frac{1}{2} \phi \left(r_{i+\frac{1}{2}}^- \right) (n_{i+1} - n_{i+2}) \right), & \text{for } W_{i+\frac{1}{2}} < 0 \end{cases}$$

The Koren flux limiter is defined as:

$$\phi(r) = \max\left(0, \min\left(2r, \min\left(\frac{1}{3} + \frac{2r}{3}, 2\right)\right)\right)$$

The upwind ratio of consecutive solution gradients:

$$r_{i+\frac{1}{2}}^{+} = \frac{n_{i+1} - n_i + \epsilon}{n_i - n_{i-1} + \epsilon}$$

$$r_{i+\frac{1}{2}}^{-} = \frac{n_i - n_{i+1} + \epsilon}{n_{i+1} - n_{i+2} + \epsilon}$$

Time integration

- Time integration is performed by employing the 2nd order Runge-Kutta method.
- Electric potential of space charge, resulting electric field, and transport coefficients are calculated at each stage of the Runge-Kutta method
- Time step restrictions:
- o The CFL condition:

$$\Delta t \left(\sum_{i=1}^{N_{dim}} \frac{|W_i|}{\Delta x} + \frac{2N_{dim}D}{\Delta x^2} \right) < \text{CFL}, CFL=0.4$$

The dielectric relaxation time:

$$\Delta t < \frac{\varepsilon_0}{q_e \mu_e n_e}, \quad \mu_e = \frac{|W(|E_{total}|)|}{|E_{total}|}$$

The AMReX library

- An open-source C++ library for massively parallel, blockstructured adaptive mesh refinement (AMR) applications
- Has inbuild geometric multigrid solvers
- Has many inbuilt classes which enable a convenient implementation of both grid and particle data
- Allows both MPI and OpenMP parallelization, as well as parallelization on graphic processing units

Adaptive mesh refinement

- Adaptive mesh refinement is used when high precision is required only in the subset of the calculation domain
- Using a uniform mesh with high resolution is impractical and inefficient under such circumstances
- The streamer dynamics is determined by the electron dynamics in the narrow region at the streamer front
- Thus, a high-resolution mesh is required at the streamer front, while a coarser mesh can be used in the streamer channel, and in the other parts of the domain, which are far from the streamer front

Refinement criteria

■ The refinement criterion due to ionization frequency:

$$\alpha_{eff}(c_1|E_{total}|)\Delta x > c_0, \qquad \alpha_{eff} = \alpha - \eta \qquad c_1 = 1.2, \quad c_0 = 0.8$$

The refinement criterion due to number density of space charges:

$$\frac{\Delta x^2 |\rho|}{\epsilon_0} > c_2, \quad c_2 = 0.1V$$

■ The de-refinement criterion:

$$\alpha_{eff}(c_1|\boldsymbol{E}_{total}|)\Delta x < 0.1 \land \Delta x < 30\mu m$$

Solving elliptic equations

- The Poisson equation and the Helmholtz equation are elliptic equations
- Their solution in each part of the domain depends on the values of the right-hand side in the entire domain
- Geometric multigrid method is very efficient in solving elliptic equations
- The method consists of applying a simple iterative procedure (like Gauss-Seidel or Jacobi) across a hierarchy of grids with reducing resolution
- AMReX includes inbuilt geometric multigrid solvers for the Poisson equation and the Helmholtz equations.
- These multigrid solvers can be easily applied across the entire hierarchy of adaptive mesh refinement levels

Results: Simulation conditions

- Axisymmetric model
- Domain size: [0 mm, 16 mm] along both coordinates
- Number of points along each axis at the coarsest level: 64
- Boundary conditions:
- For the number density of electrons: Zero Neumann conditions at all boundaries
- For the electric potential of space charge and photoionization source terms: Zero Neuman conditions at boundaries perpendicular to the radial axis and zero Dirichlet conditions at boundaries perpendicular to the axial coordinate
- Results of our AMReX-streamer code are compared to the results of the open-source Afivo-streamer code

Two-headed streamer in the Titan mixture

- The Titan mixture: 98.4% N₂, 1.6% CH₄
- The number of AMR levels 7 (from 0 to 6)
- No photoionization
- The applied reduced electric field: $E/N_0 = 147 \text{ Td}$, $1Td = 10^{-21} \text{ V/m}$, $N_0 = 2.50475764 \cdot 10^{25} \text{ m}^{-3}$.
- The initial condition: Neutral Gaussian given by:

$$n_e = n_p = n_{background} + n_0 \cdot e^{-\frac{(z-z_0)^2 + r^2}{\sigma^2}}, \quad n_n = 0m^{-3}$$

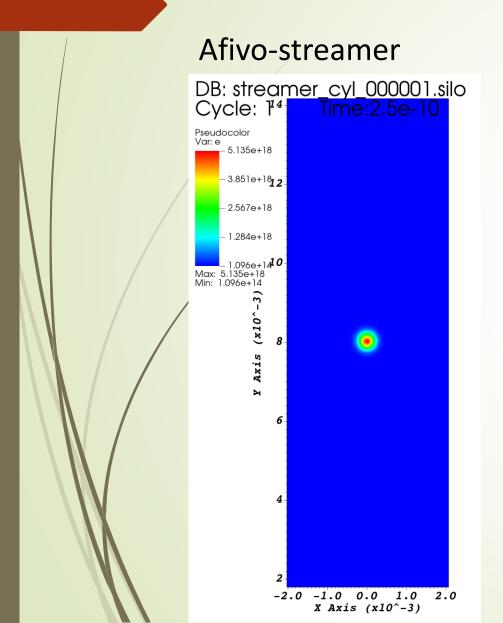
$$n_{background} = 10^{14} m^{-3}$$

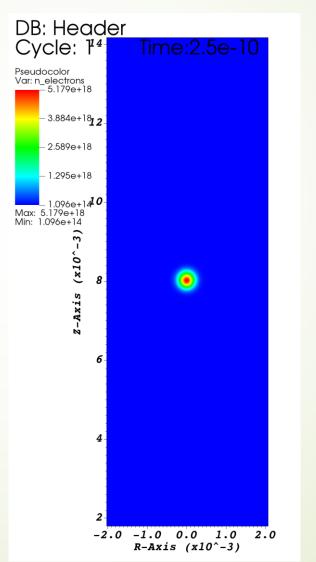
$$n_0 = 5 \cdot 10^{18} m^{-3}$$

$$z_0 = 8mm$$

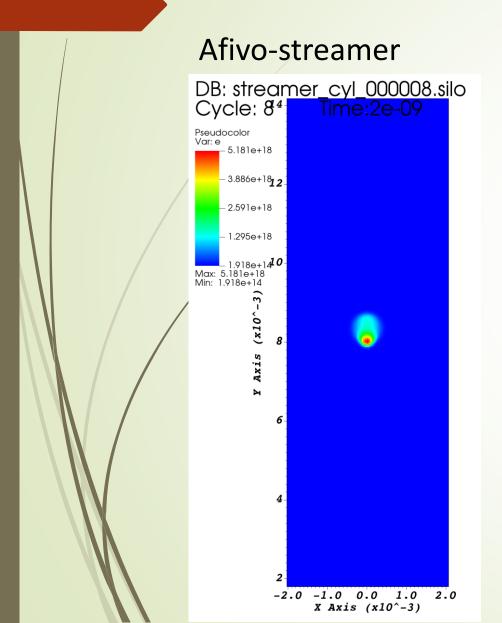
$$\sigma = 4 \cdot 10^{-4} m$$

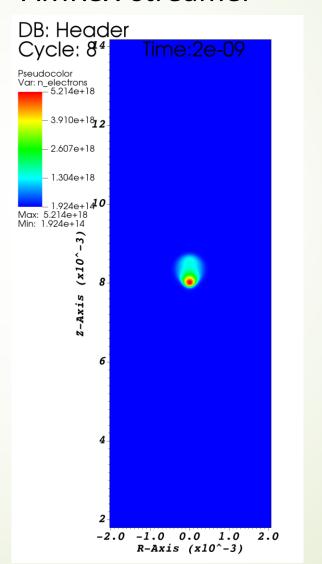
Number density of electrons at t = 0.25 ns



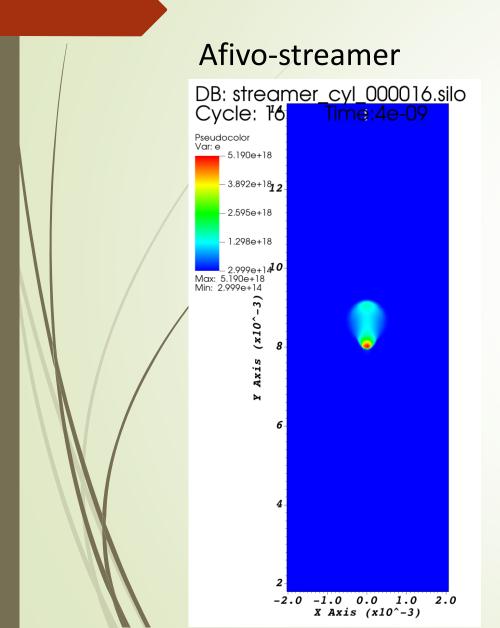


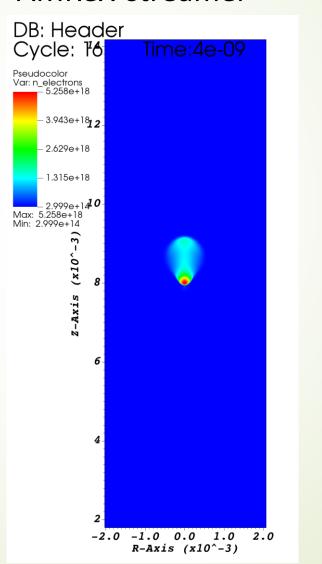
Number density of electrons at t = 2 ns



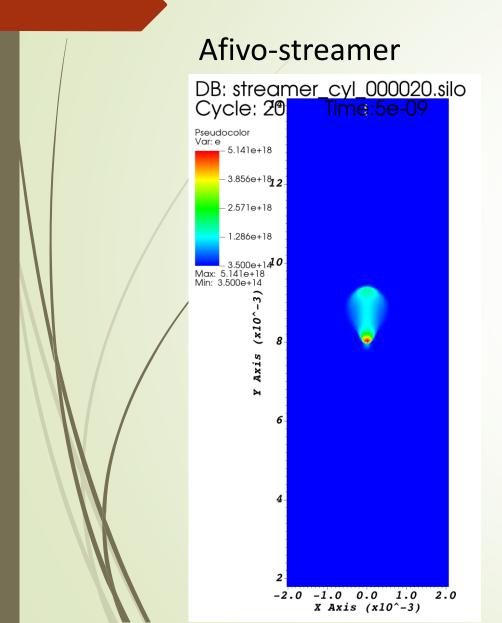


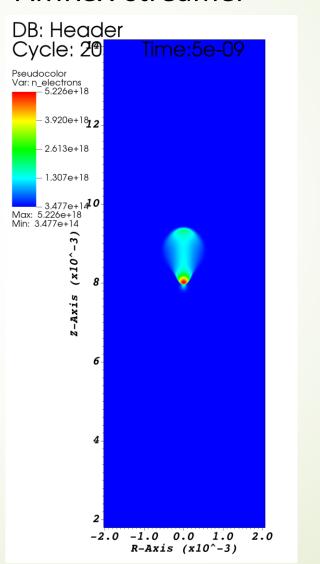
Number density of electrons at t = 4 ns



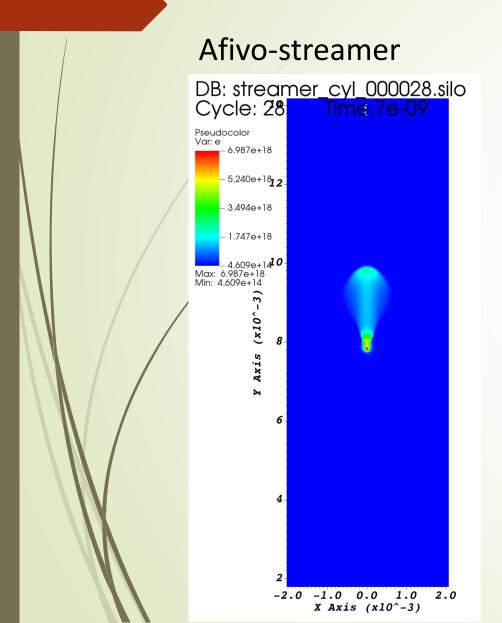


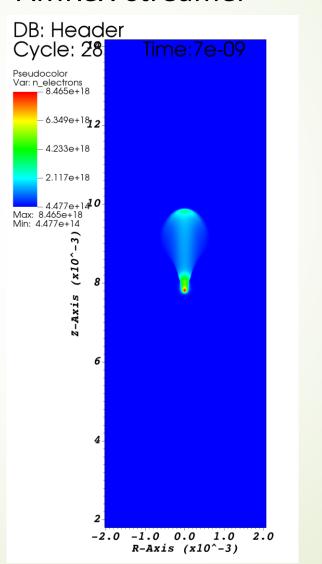
Number density of electrons at t = 5 ns



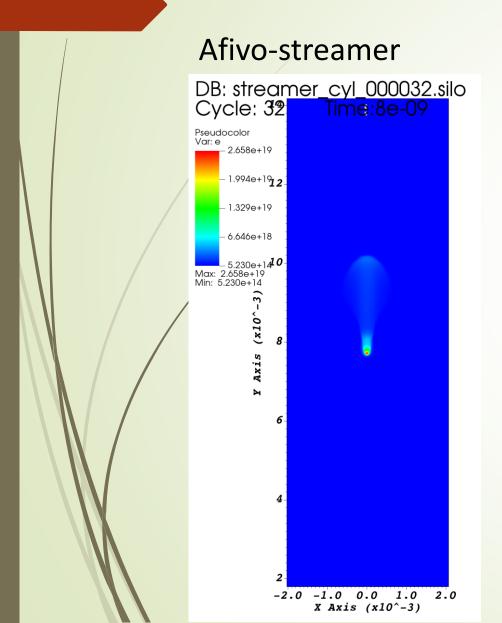


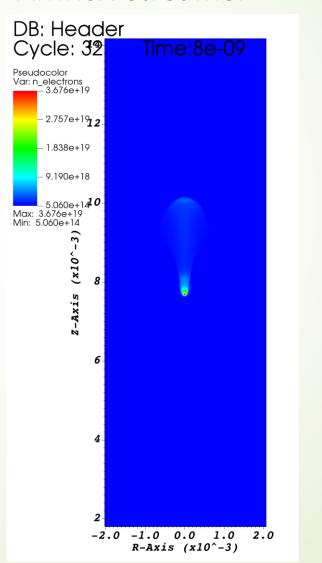
Number density of electrons at t = 7 ns



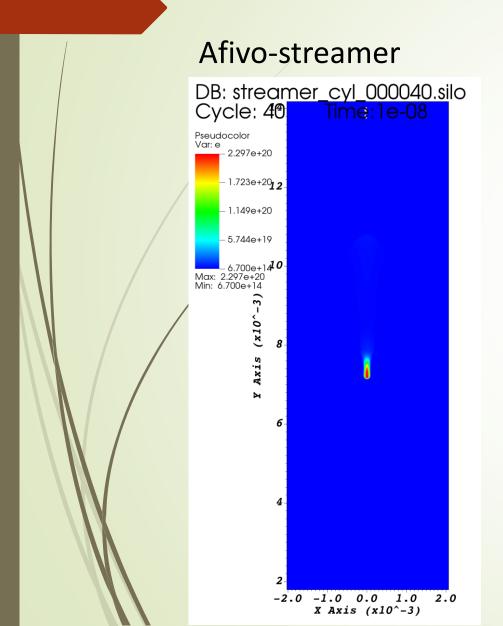


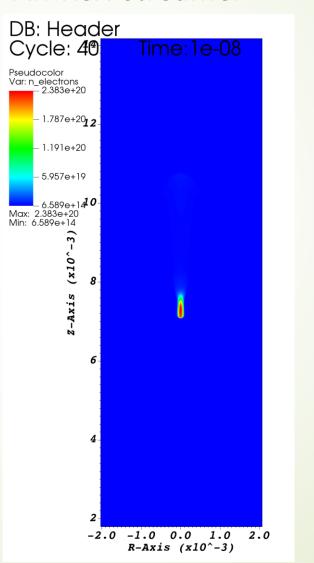
Number density of electrons at t = 8 ns



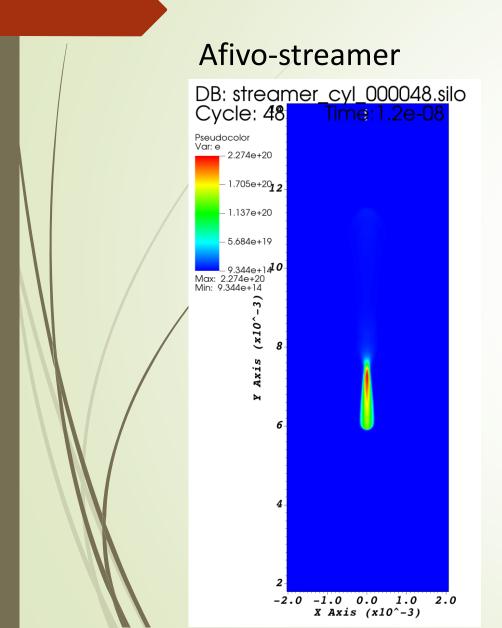


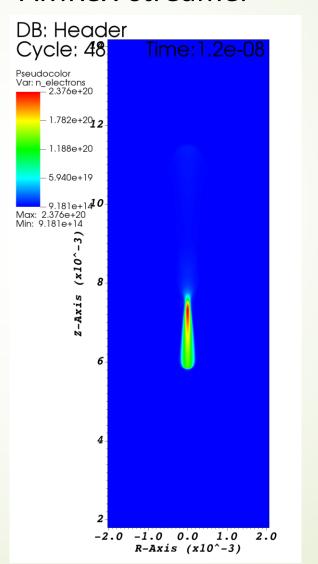
Number density of electrons at t = 10 ns



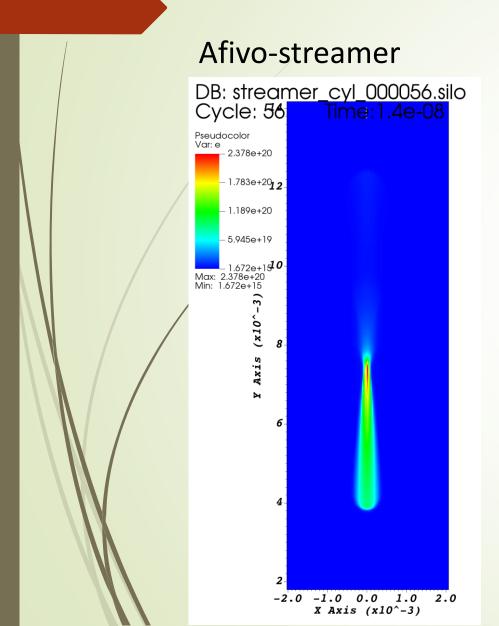


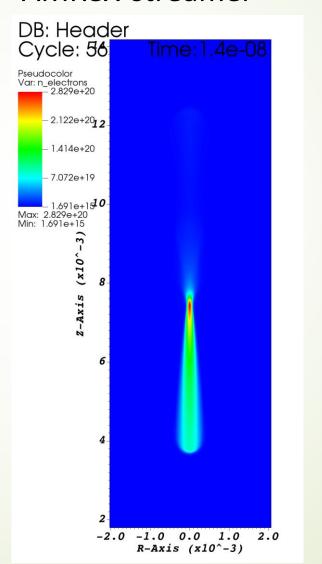
Number density of electrons at t = 12 ns



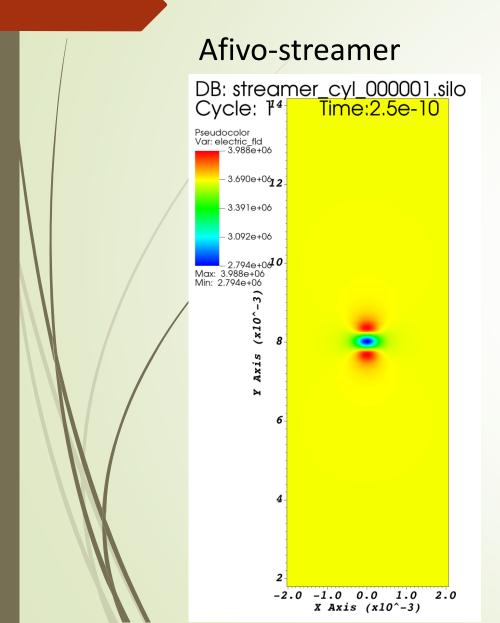


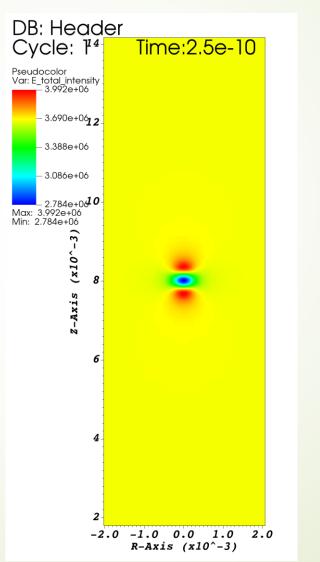
Number density of electrons at t = 14 ns



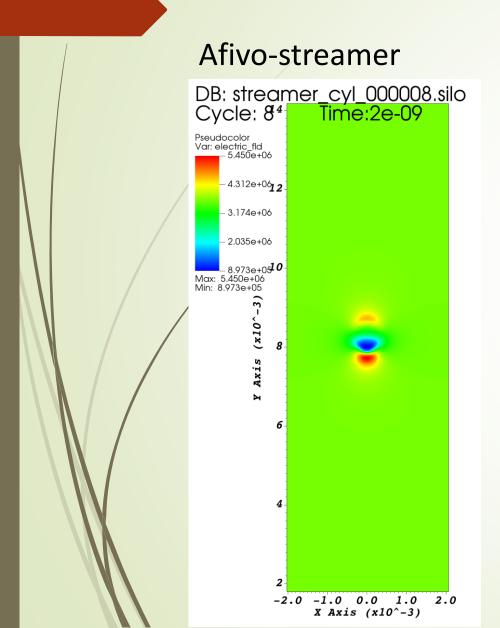


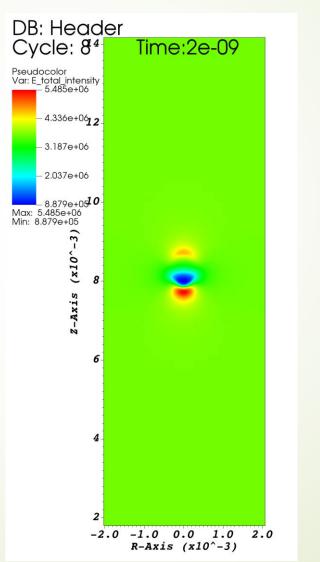
Electric field intensity at t = 0.25 ns



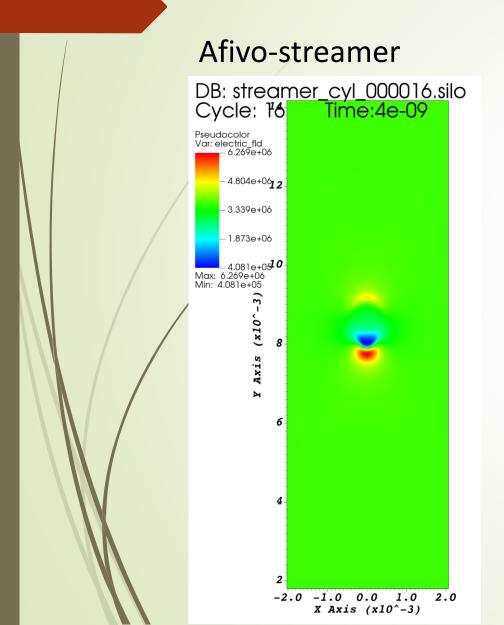


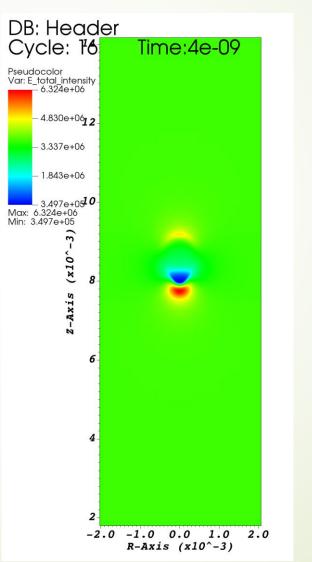
Electric field intensity at t = 2 ns



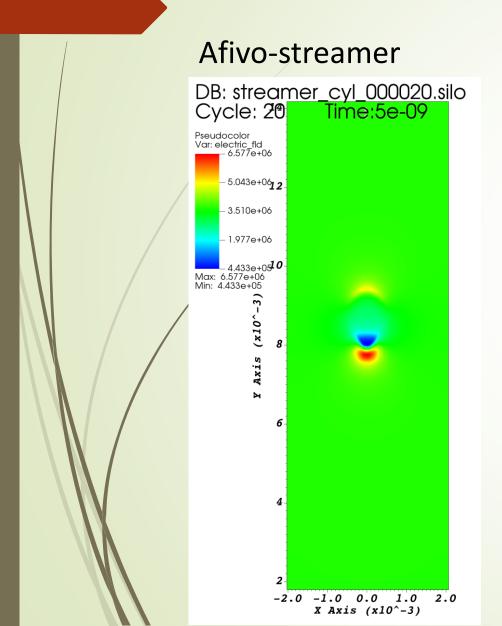


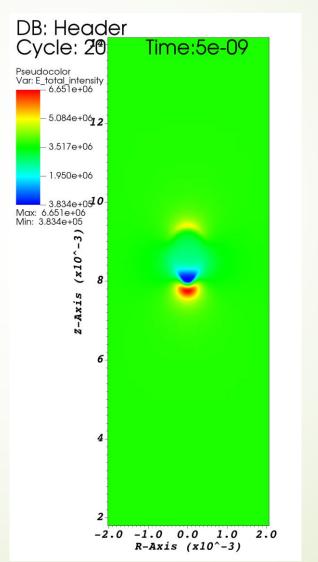
Electric field intensity at t = 4 ns



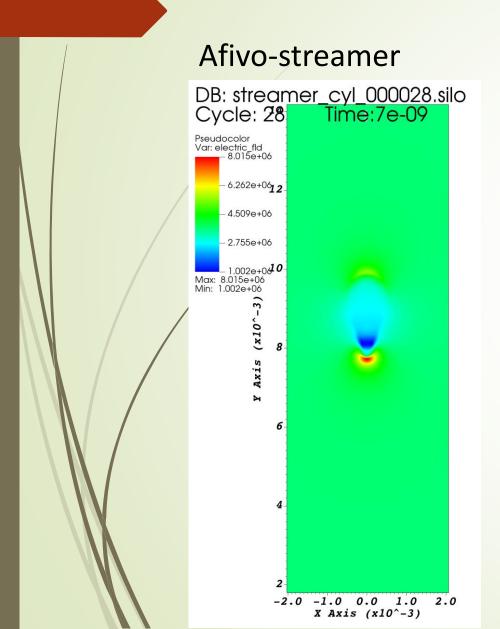


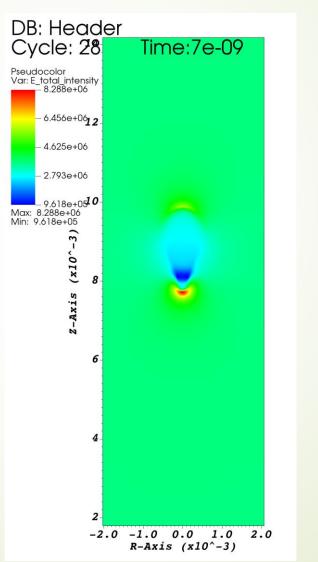
Electric field intensity at t = 5 ns



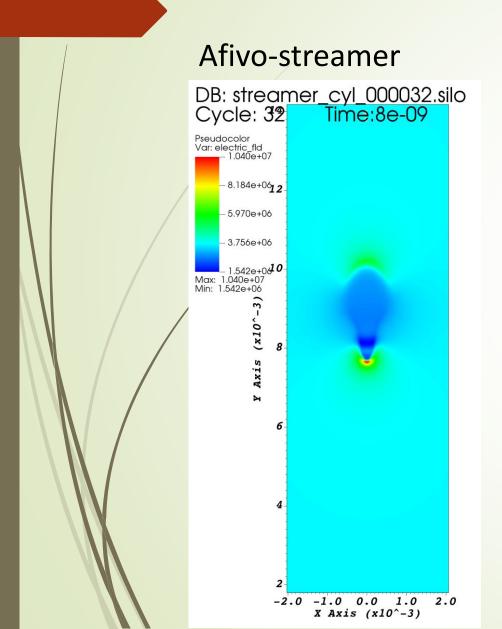


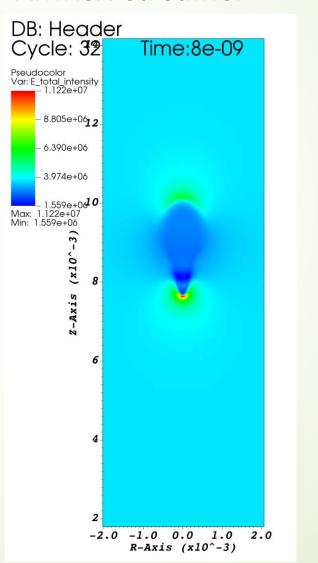
Electric field intensity at t = 7 ns



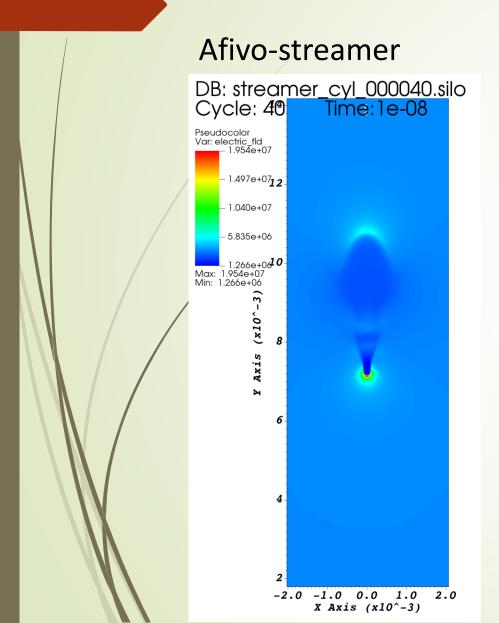


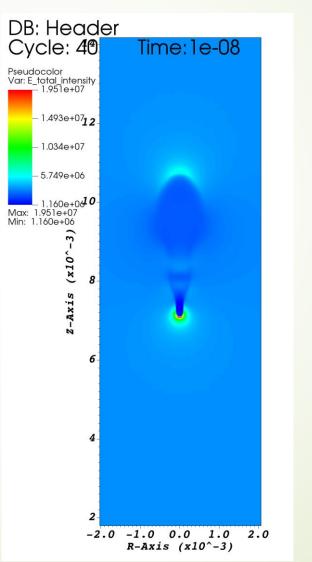
Electric field intensity at t = 8 ns



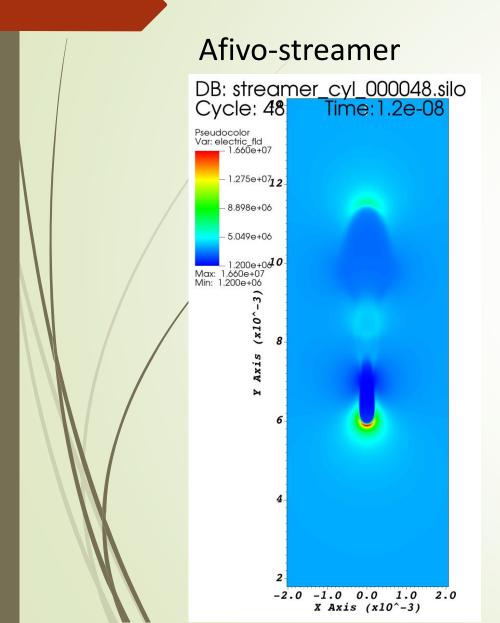


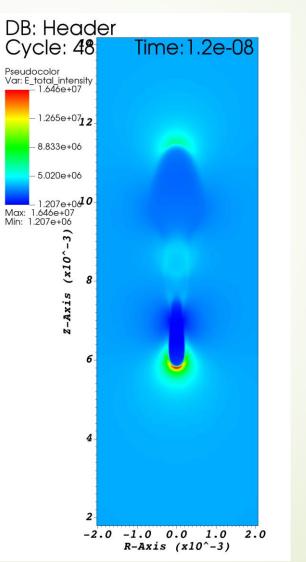
Electric field intensity at t = 10 ns



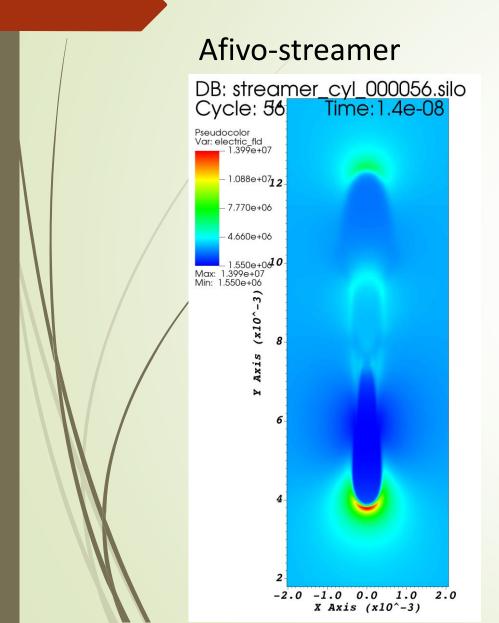


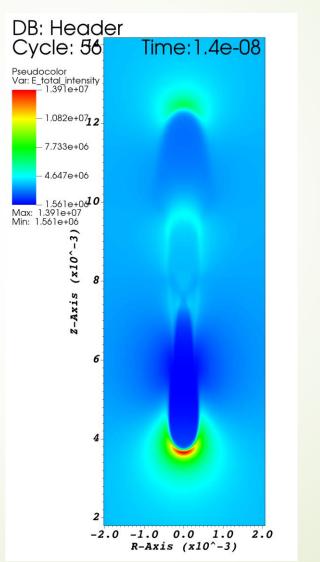
Electric field intensity at t = 12 ns





Electric field intensity at t = 14 ns



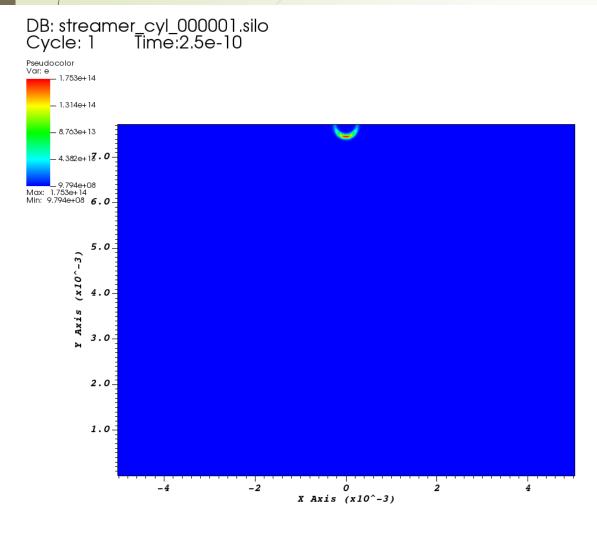


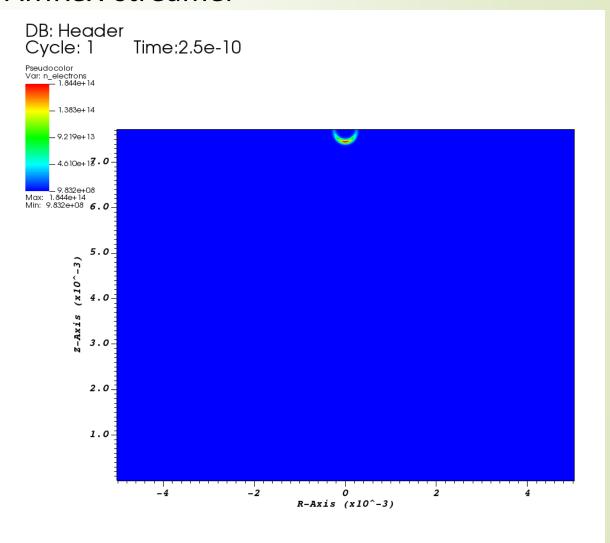
Positive streamer in the air

- The dry air: $80.0\% N_2$, $20.0\% O_2$.
- The number of AMR levels 7 (from 0 to 6)
- Photoionization parametrization: Bourdon 3 term
- The applied reduced electric field: $E/N_0 = 80 \text{ Td}$
- The initial condition: Positive line segment with a length of 0.64 mm and a width of 0.4 mm, centered at the center of the z coordinate range
- Background ionization: 10⁹m⁻³.
- Neutral gas density: 2.50475764 ⋅ 10²⁵ m⁻³.

Number density of electrons at t = 0.25 ns

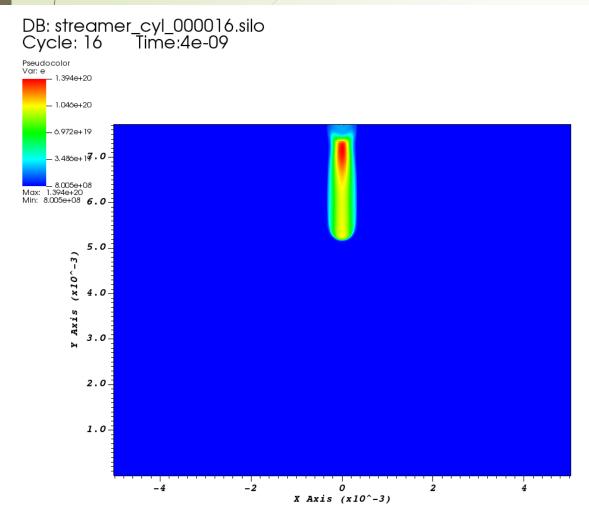
Afivo-streamer

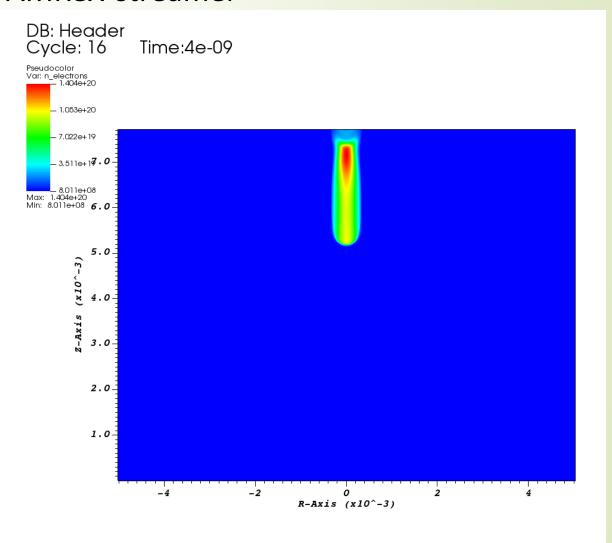




Number density of electrons at t = 4 ns

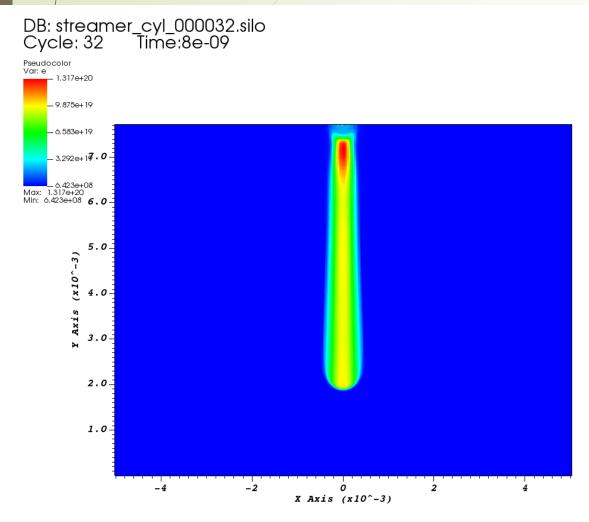
Afivo-streamer

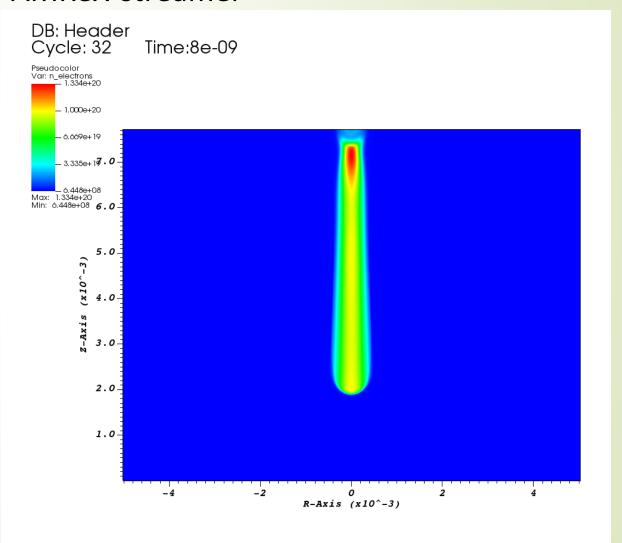




Number density of electrons at t = 8 ns

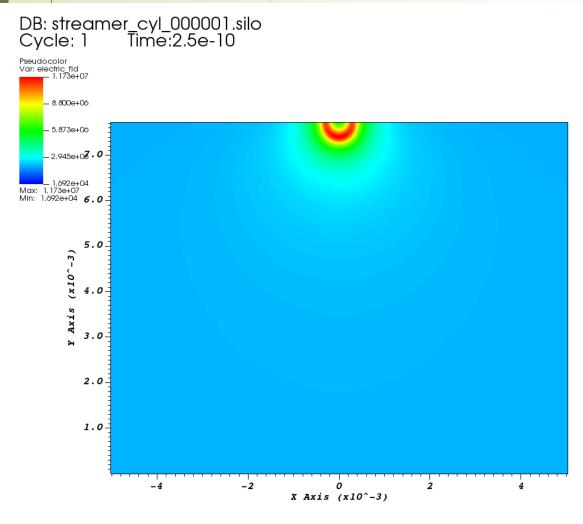
Afivo-streamer

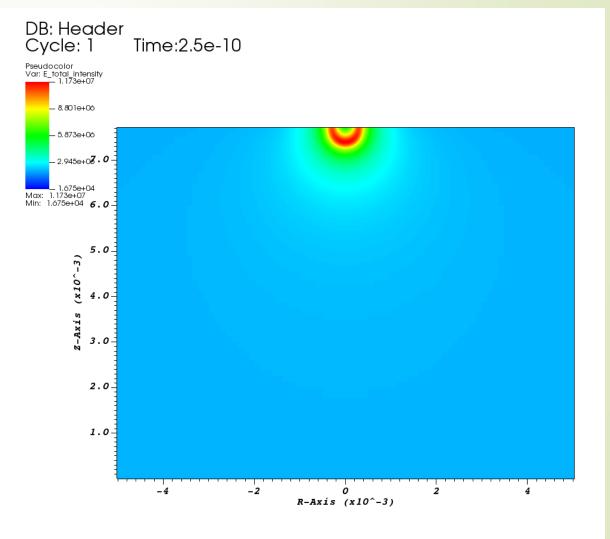




Electric field intensity at t = 0.25 ns

Afivo-streamer

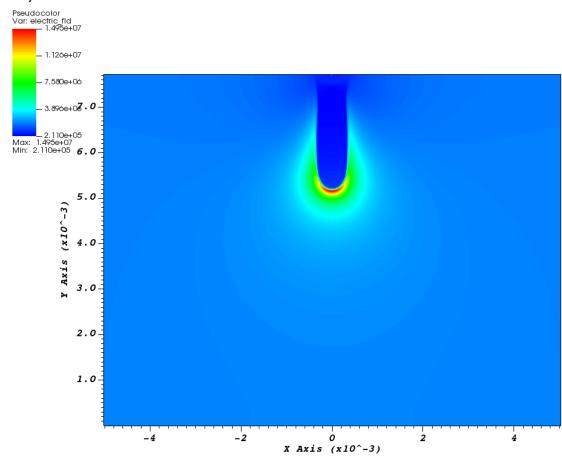


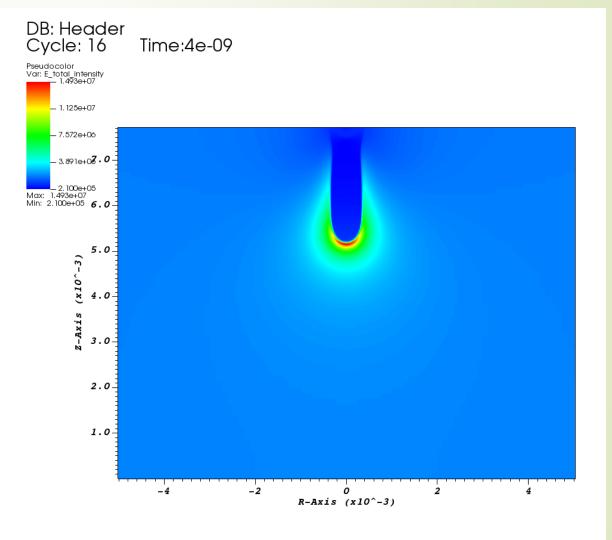


Number density of electrons at t = 4 ns

Afivo-streamer



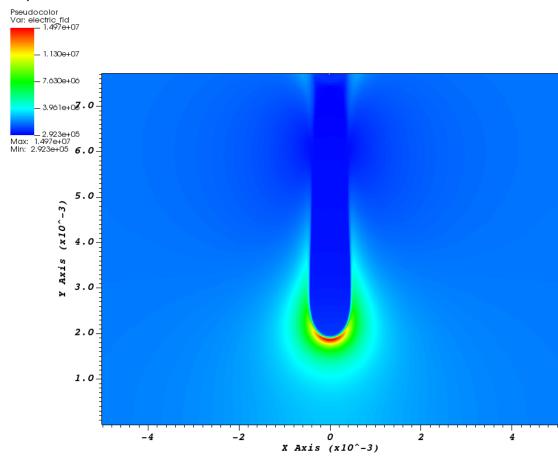


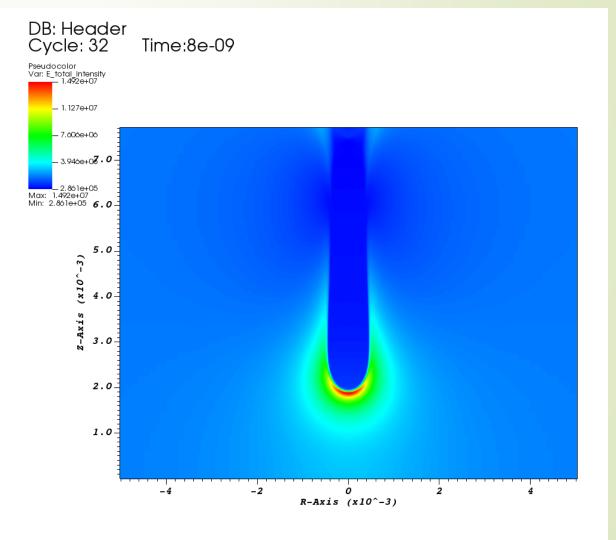


Number density of electrons at t = 8 ns

Afivo-streamer







Conclusions

- We have implemented an axisymmetric fluid model in the AMReX library
- Our model is based on the first-order fluid model with local field approximation
- Photoionization is implemented by solving a system of Helmholtz equations
- The good agreement between the results of our program and the Afivo-streamer open-source program confirms the validity of our code.

Acknowledgements

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- The authors acknowledge support from the MESTD of Serbia.

THANK YOU